1	Phase change dynamics in a cylinder containing hybrid
2	nanofluid and phase change material subjected to a
3	rotating inner disk
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15 Abstract

In this numerical study, the phase change dynamics of a 3D cylinder con-16 taining hybrid nanofluid and phase change material (PCM) is investigated 17 with a finite element solver. The PCM consists of spherical encapsulated 18 paraffin wax, and the flow is under the forced convection regime. The dy-19 namic features of the phase change process are studied for different values of 20 the Reynolds number (between Re=100 and 300), the rotational Reynolds 21 number of the inner disk (Rew=0 and 300), and the size of the rotating disk 22 (length between 0.1L and 0.55L; height between 0.001H2 and 0.4H2). The 23 flow dynamics and separated flow regions are found to be greatly influenced 24 by the rotational speed and size of the inner disk. As Re is increased, the 25 difference between the transition times at different rotational disk speeds de-26 creases. At Re=100, a 21% reduction in the phase transition time is observed 27

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when the inner disk rotates at the highest speed as compared to the motion-28 less case. Up to a 26% variation in the phase transition time occurs when 29 the size of the inner rotating disk is varied. A 5 input-1 output feed-forward 30 artificial neural network is applied to achieve fast and reliable predictions of 31 the phase change dynamics. This study shows that introducing rotational 32 effects can have a profound effect on the phase change dynamics of a hybrid 33 nanofluid system containing phase change material. 34

Keywords: rotational surface, phase change, CFD, finite element method, 35

artificial neural network, hybrid particles 36

Nomenclature

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k thermal conductivity	T_m melting temperature
H cylinder height	u, v, w velocity components
L cylinder length	r, z cylindrical coordinates
L_f latent heat of fusion	Greek Characters ν kinematic viscosity
n unit normal vector	ρ density of the fluid
Nu Nusselt number	ϕ solid volume fraction
p pressure $L_{1,2}$ disk radius	ω rotational speed
Pr Prandtl number	Subscripts
Re Reynolds number	$\begin{array}{c} c & \text{cold} \\ h & \text{hot} \end{array}$
Rew rotational Reynolds number	<i>m</i> average
t time T temperature	<i>nf</i> nanofluid
	p solid particle

38 1. Introduction

Energy applications involving phase change material (PCM) equipped 39 products have been gaining attention recently due to their performance en-40 hancement potential. Thermal energy (TE) storage and management are 41 two major areas that are applicable to various energy systems such as in 42 solar power, refrigeration, electronic cooling, waste heat recovery, building 43 energy, agriculture and many others. Due to the low cost and favorable 44 heat transfer (HT) characteristics, packed bed latent heat storage (HS) is 45 considered a promising technology. Many methods can improve the per-46 formance features of PCM embedded TE systems. In the review of Khan 47 et al. [1], different methods for influencing the thermo-physical properties 48 of PCMs were presented such as using fins, high conductivity additives and 49 various arrangement of PCMs in TE systems. Thermal utilization of PCM 50 was explored by considering different techniques and accounting for stabil-51 ity issues. Sharma et al. [2] reviewed PCM installed TE systems in diverse 52 applications such as in waste heat, building and solar air heaters. Enhanc-53 ing the thermal performance of PCMs by using internal fins, metallic foams 54 and nano particles was considered in the review by Sahoo et al. [3]. Heat 55 sinks with PCM enhancement techniques were considered for constant and 56 variable loads while metallic foams offered several advantageous. Fins with 57 PCM have been considered in heat sinks [4, 5], photovoltaic panels [6, 7] and 58 building applications [8, 9]. 59

Recently, nano-sized particles have been used in HT fluids for TE systems to improve performance. Many nanofluid types have been evaluated and experimental correlations for the effective thermophysical properties have been

derived [10]. Many advanced simulation methods have been developed and 63 tested for accurate modeling of TE systems using nanofluids [11-15]. The 64 utilization of nanotechnology in PCM embedded systems has been consid-65 ered in various TE systems [16]. In the review of Rostami et al. [17], PCM 66 characteristics including nano-PCMs on TE storage and natural convection 67 were critically discussed. Through numerical simulations, Khodadadi and 68 Hosseinizadeh [18] showed performance improvement for freezing of water 69 by using CuO nanoparticles in a cavity with numerical simulations. Other 70 studies that have considered the use of nanoparticles with PCM can be found 71 in Refs. [19–27]. 72

The performance of PCM embedded systems can be enhanced by chang-73 ing the geometrical features along with the thermophysical properties of the 74 HT fluid. In TE systems with embedded latent heat storage units with 75 encapsulated PCM, dynamic walls can be implemented to improve the per-76 formance. Rotating surfaces have been considered in many TE systems espe-77 cially for convective HT applications. Although the thermal system perfor-78 mance can be improved via passive methods such as changing the geometry 79 of the TE systems and installing static elements such as internal fins, metal 80 foams, the use of dynamic walls or objects opens up new opportunities for 81 further improvements. A rotating type object has been shown to affect the 82 convective HT thermal performance via adjustments in the rotational speed. 83 size and location of the object [28, 29]. Certain locations and rotational 84 direction of the cylinder have been found to assist the convection in TE sys-85 tems. The combined effects of using nanofluids and surface rotations have 86 been considered in many studies for HT improvement in TE systems [30, 31]. 87

The present work considers the effects of using a rotating inner disk on the 88 performance of a PCM embedded thermo-fluid system. The effects of rotation 89 of the inner disk and of its geometrical parameters on the dynamic features 90 of the phase change process are numerically assessed. The effects of fluid 91 velocity and its interaction with the rotation of the inner disk surface on the 92 flow features are analyzed. The numerical simulations are validated against 93 an experimental study of the phase change process. With the diverse use of 94 PCMs in TE systems such as in electronic cooling, solar energy applications, 95 heat exchangers, the thermal management of the phase change process via a 96 rotating inner disk and nanofluid is novel and may be used to enhance the 97 performance of PCM embedded energy systems. 98

99 2. Computational study

100 2.1. System configuration

The characteristics of a PCM-equipped cylinder under the effects of a 101 rotating inner disk is analyzed. The height and radius of the PCM included 102 region is hpcm and L. The rotating disk has inner and outer radius of L1d 103 and L2d while heights are H1d and H2d. The rotational speed (Rs) of the 104 inner disk is ω . Here, H1 and H2 denote the distance of the PCM region 105 from the outlet and inlet, respectively. A hybrid nanofluid (with Ag and 106 MgO nanoparticles at solid volume fraction of 2%) enters the cylinder with 107 a velocity of ug and a temperature of Tg. Experimental correlations for the 108 effective nanofluid properties were used as given in [32]. Spherical shaped 109 encapsulated paraffin wax is considered as the PCM with a radius of 30 mm. 110 Table 1 lists the thermophysical properties. 111

112 2.2. Governing equations and boundary conditions

A single phase model of nanofluid is used with Newtonian and incompressible fluid assumptions. Effects such as free convection, thermal radiation and viscous dissipation are ignored.

A hybrid nanofluid containing binary particles of Ag and MgO was used. 116 The potential of using hybrid nanofluids has been shown in various studies 117 of thermal engineering. They are preferred for their synergistic effects, cost 118 and advantages of one or more types of nanoparticles in the base fluid [33-119 36]. In studies with nanofluid, an accurate description of the thermophysical 120 properties is important. The experimental data fit from Ref. [32] was used to 121 derive correlations for the thermal conductivity (k_{nf}) and viscosity (μ_{nf}) of 122 the hybrid nanofluid. They are defined as [32]: 123

$$k_{nf} = \left(\frac{0.1747 \times 10^5 + \phi}{0.1747 \times 10^5 - 0.1498 \times 10^6 \phi + 0.1117 \times 10^7 \phi^2 + 0.1997 \times 10^8 \phi^3}\right) k_f,$$
(1)
$$\mu_{nf} = \left(1 + 32.795 \phi - 7214 \phi^2 + 714600 \phi^3 - 0.1941 \times 10^8 \phi^4\right) \mu_f.$$
(2)

where ϕ is the total concentration of nanoparticles and is taken as 2%. In the fluid domain, the conservation equations are [37]:

$$\nabla \mathbf{.u} = \mathbf{0} \tag{3}$$

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$$\rho\left(\mathbf{u}.\nabla\right)\mathbf{u} = \nabla\left[-p\mathbf{I} + \mathbf{K}\right] + F \tag{4}$$

$$\mathbf{K} = \mu \left(\nabla \mathbf{u} + \left(\nabla \mathbf{u} \right)^T \right) \tag{4}$$

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$$\rho C_p \frac{\partial T}{\partial t} + \rho C_p \mathbf{u} \cdot \nabla T = \nabla \cdot (k \nabla T)$$
(5)

¹²⁹ For the PCM region:

$$\nabla . \mathbf{u} = \mathbf{0} \tag{6}$$

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$$\frac{1}{\varepsilon_p} \rho \left(\mathbf{u} \cdot \nabla \right) \mathbf{u} \frac{1}{\varepsilon_p} = \nabla \cdot \left[-p\mathbf{I} + \mathbf{K} \right] - \left(\mu \kappa^{-1} + \beta_p \rho |\mathbf{u}| \right) \mathbf{u} + F$$

$$\mathbf{K} = \mu \frac{1}{\varepsilon_p} \left(\nabla \mathbf{u} + \left(\nabla \mathbf{u} \right)^T \right)$$
(7)

¹³¹ and with the Kozeny-Carman permeability given as:

$$\kappa = \frac{d_p^2}{180} \frac{\varepsilon_p^3}{(1 - \varepsilon_p)^2} \tag{8}$$

¹³² A phase change function α is defined with a value of 0 for $T < (T_m - \Delta T_m/2)$ ¹³³ and 1 for $T > (T_m + \Delta T_m/2)$. The energy equation is given as:

$$\rho C_p \frac{\partial T}{\partial t} + \rho C_p \mathbf{u} \cdot \nabla T = \nabla \cdot \left(k \nabla T \right).$$
(9)

¹³⁴ with the thermo-physical properties as:

$$\theta = 1 - \alpha, \quad \rho = \theta \rho_{f1} + (1 - \theta) \rho_{f2},$$

$$C_p = \frac{1}{\rho} \left(\theta \rho_{f1} C_{p,f1} + (1 - \theta) \rho_{f2} C_{p,f2} \right) + L \frac{\partial \alpha_m}{\partial T},$$

$$k = \theta k_{f1} + (1 - \theta) k_{f2}, \quad \alpha_m = \frac{1}{2} \frac{(1 - \theta) \rho_{f2} - \theta \rho_{f1}}{\theta \rho_{f1} + (1 - \theta) \rho_{f2}}.$$
(10)

In the above derivation, f1, f2 and L are the phases and latent heat of fusion. In between the phases, non-equilibrium HT in porous media interface is considered. Here, two equations for the phases and additional source terms are given [38, 39]:

$$\theta_p \rho_s C_{p,s} \frac{T_s}{\partial t} + \nabla \mathbf{q_s} = q_{sf} \left(T_f - T_s \right) + \theta_p Q_s, \tag{11}$$

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$$\mathbf{q}_{\mathbf{s}} = -\theta_p k_s \nabla T_s,\tag{12}$$

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$$(1 - \theta_p)\rho_f C_{p,f} \frac{T_f}{\partial t} + (1 - \theta_p)\rho_f C_{p,f} \mathbf{u_f} \cdot \nabla T_f +$$
(13)

$$+\nabla \mathbf{q_f} = q_{sf}(T_s - T_f) + (1 - \theta_p)Q_f,$$

$$\mathbf{q}_{\mathbf{f}} = -(1 - \theta_p)k_f \nabla T_f. \tag{14}$$

In the above equations, $\mathbf{q_f}$, $\mathbf{q_s}$, θ_p denote the conductive heat fluxes of the solid and fluid and porous media SVF. The interstitial convective HT coefficient is given by the term q_{sf} while Q_s and Q_f are the solid and fluid heat sources, respectively. For a spherical pellet bed [38, 39]:

$$q_{sf} = a_{sf} h_{sf} \tag{15}$$

¹⁴⁶ The interstitial HT coefficient is given by [38, 39]:

$$\frac{1}{h_{sf}} = \frac{2r_p}{k_f \text{Nu}} + \frac{2r_p}{\beta k_s}.$$
(16)

¹⁴⁷ The value of β is taken to be 10 for spherical particles. The fluid-to-solid Nu ¹⁴⁸ number is: [40]:

$$Nu = 2 + 1.1 Pr^{1/3} Re_p^{0.6}.$$
 (17)

¹⁴⁹ The Prandtl number and Reynolds number of the particles are:

$$\Pr = \frac{\mu C_{p,f}}{k_f}, \quad \operatorname{Re}_p = \frac{2r_p \rho_f |u_f|}{\mu}.$$
(18)

The relevant non-dimensional parameters are the Reynolds number (Re) for the fluid and the rotational Reynolds number (Rew) for the rotating inner disk, both of which are varied in our simulations. They are given as:

$$Re = \frac{u_g L\rho}{\mu}, \quad Rew = \frac{\omega (L1d)^2 \rho}{\mu}$$
(19)

The Re number and rotational Re of the inner disk are the varied during the simulation. The fluid velocity at the inlet is ug and the temperature is Tg=336 K. The model is axis-symmetrical with $\frac{\partial T}{\partial r} = 0$. The walls of the cylinder are adiabatic with no-slip boundary conditions (BCs) as, u = w = 0, $\frac{\partial T}{\partial n} = 0$. A pressure outlet is used at the exit. The inner disk is rotating with Rs of ω . The velocity of the tangentially moving wall is 0 while the angular component of the moving wall is ωr . The rotating disk is also considered adiabatic such that $\frac{\partial T}{\partial n} = 0$ with an initial temperature of 303 K.

161 2.3. Solution method and code validations

To solve the above GEs, the GWR-finite element method is used. The field variable (Ψ) appropriations are performed as:

$$\psi = \sum_{n=1}^{N^s} \Phi_n^s \Psi_n \tag{20}$$

where Φ^s is the shape function and Ψ is the element nodal value. Different 164 ordered Lagrange FEM is considered. The time dependent part is treated 165 by utilizing a second order backward differentiation scheme. A time step 166 size of 0.1 min is used. Time step independence of the solution is also as-167 sured. A commercial computational fluid dynamics code Comsol was used 168 with built-in modules for multi-physics simulations [39]. Tests for grid inde-169 pendence were conducted and the results are given in Figure 2 (a) for two 170 different values of Rs of the inner disk. Grid system G4 with 63412 elements 171 is selected. At the interface regions and near the walls, the mesh is finer 172 as shown in Figure 2 (b). Validation is conducted by using results for the 173 phase change process and impacts of rotation on the convective HT. First, 174 results available in Ref. [41] are used. In the study, experimental analysis for 175 PC process in a differentially heated cavity was performed. The role of free 176 convection for PC of pure metal was explored while the two vertical walls of 177 the rectangular test section were maintained at different temperatures. The 178

amount of solidified volume is a time dependent quantity and it depends on the solidified volume (V), Rayleigh number (Ra), total volume (V_0) , aspect ratio (AT) and dimensionless time (τ) as:

$$\frac{V}{V_0} = 2.91 \tau^{0.53} \text{Ra}^{-0.05} A^{-0.36}$$
(21)

Figure 3 compares the solidified volume at different dimensionless times Com-182 pared with experimental data, the greatest difference observed as -9.3% at 183 $\tau^{0.53} \text{Ra}^{-0.05} A^{-0.36} = 0.114$. In another study, Roslan et al. [42] explored the 184 effects of rotation on convective HT by using an inner cylinder in square cav-185 ity. They considered a differentially heated enclosure while rotational speed 186 and size effects of the inner cylinder on the convective HT features were ex-187 plored using FEM. Comparison of the average Nu is given in Figure 5 for two 188 different configurations (cylinder size and Rs). In both cases, the difference 189 is limited to under 3%. The results reinforce the capability of the code in 190 simulating the PC process and the rotational effects of the inner surfaces on 191 convection. 192

193 2.4. Performance estimation with artificial neural network

Artificial neural network (ANN) modeling is widely used in energy systems and thermal engineering for flow control, performance prediction and system identification. Many different soft computing methods and ANN techniques have been demonstrated for accurate performance predictions of energy systems [43–49]. Many steps involved in soft computing methods are similar such as specification of the input-output data set, activation functions, learning algorithms and validation tests. The Levenberg-Marquardt

(LM) techniques with backpropogation (BP) is chosen here as learning algo-201 rithm [50]. A mathematical model for the input-output relation of the data 202 set can be derived with an ANN. The network structure consists of multiple 203 different layers (input, hidden and output layers) and different numbers of 204 neurons are used in the hidden layers. They can be adjusted along with the 205 network parameters during the training phase. The network performance at 206 each iteration step is checked by comparing the network output and data 207 from simulations or experiments while the weights of the ANN are updated. 208 A learning algorithm with backpropagation is used for adjusting the weights 209 of ANN. Each neuron output is given as [47]: 210

$$Y = G\left(\sum_{i=1}^{M} X_i W_i + S\right) \tag{22}$$

where Y, G and M represent the output data, activation function and data number while X, W and S are the input, weight and bias term, respectively. Different criteria may be applied to test the performance of different ANN models. The mean square error (MSE) and coefficient of determination (\mathbb{R}^2) are commonly used [51]:

$$MSE = \frac{1}{M} \sum_{i=1}^{M} \left(y_i^{CFD} - y_i^N \right)^2,$$
(23)

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$$R^{2} = 1 - \frac{\sum_{i=1}^{M} (y_{i}^{CFD} - y^{N})^{2}}{\sum_{i=1}^{M} (y_{i}^{CFD} - \bar{y}^{N})^{2}}.$$
(24)

217 3. Results and discussion

The effects of using a rotating object in a PCM filled cylinder on the performance enhancement are numerically assessed. The Rs of the disk is ω

and is varied during the simulation. The dynamic characteristics of the phase 220 change process are studied for various values of Re (between 100 and 300), 221 Rew of the disk (0 and 300), and size of the rotating disk (length between 222 0.1L and 0.55L; height between 0.001H2 and 0.4H2). The time evolution of 223 the liquid fraction (Lf) and phase completion time are analyzed. ANN-based 224 estimation is performed to determine the time dependent characteristics of 225 the liquid fraction for various values of input parameters such as Rew and 226 the size of the inner rotating disk. 227

228 3.1. CFD simulation results

Figure 5 shows the impacts of Re on the flow patterns variation at two 229 different values of Rew. When the inner disk is not rotating, the flow sep-230 arates at the edge of the disk away from the inlet and its extent increases 231 with Re. However, in the presence of rotational effects due to the inner disk 232 surfaces, the separated flow region extending through the PCM domain be-233 comes significant. As Re is increased, the size of the vortex at the interface 234 becomes reduced and the recirculation zone is suppressed. The recirculation 235 zone becomes larger at the edges of the rotating disk as Rs increases (Figure 236 6). The separated flow region moves toward the vertical wall of the cylin-237 drical container for higher Rew. Time evolutions of the PCM temperatures 238 for different Rew are presented in Figure 7. Due to the recirculation zones 239 caused by rotation of the inner disk, the PCM temperatures are significantly 240 affected in the inner part and toward the walls when the disk is rotating at 241 the highest speed. Variation of the PCM temperatures at two points (in the 242 axis- symmetric location and in the vertical wall) are shown in Figure 8 for 243 various values of Re and Rs of the inner disk. In the absence and presence of 244

rotational effects, when the fluid velocity increases, the phase change process 245 speeds up. However, the effects of Rw on the phase change differs at the mid 246 point (axis-symmetrical location) and in the wall location. At the wall loca-247 tion, phase transition becomes rapid with increasing Rew while the effects is 248 reversed at mid point location. Time evolution of the liquid fraction at two 249 different values of Rew is given in Figure 9. At t=40 min full phase transi-250 tion is achieved when rotational effects of the inner disk are considered at the 251 highest speed. Near the wall region and interior of the cylinder, rotational 252 effects on the phase transition become dominant. The dynamic characteris-253 tics of the liquid fraction (Fr) for various Re at two different values of Rew 254 of 0 and 300 are shown in Figure 10. The values of Fr approach for higher 255 Re for both values of Rew. A saturation type curve is obtained while the 256 discrepancy between different Re on Fr becomes different when rotational 257 effects are considered. Full phase transition time (tr) with different Re and 258 Rew are given in Figure 11. The effects of Rew on tr becomes dominant for 250 lower fluid velocities but at Re=300, its impacts are negligible. There is a 260 21% reduction in the full transition time at Re=100 when rotational effects 261 at the highest speed are considered as compared to a motionless disk. When 262 the effects of Rew are considered, the full phase transition time first increases 263 up to Rew=50 but decreases thereafter. This is attributed to the balance 264 between vortex formation which has negative effects on the phase change 265 process while for higher rotational speeds, the fluid velocity near the rota-266 tion surfaces becomes higher, causing a positive impact on the phase change 267 process. There is almost a 16 % reduction in tr when the configuration with 268 higher Rs is compared with the case of a motionless disk. 269

Geometrical parameters of the rotating disk have also affected the phase 270 change process. The flow pattern variations with varying length and height of 271 the rotating disk are given in Figure 12. The separated flow region becomes 272 significant and occupies a large portion of the PCM region for L2d=0.3L 273 while it is smallest for L2d=0.1L. At L=0.55L, the region approaches the 274 vertical wall of the cylinder. When the height is increased, an elongation of 275 the vortex region in the flow direction is observed. At the highest height, 276 its size is reduced in the radial direction. Varying the length and height of 277 the rotating inner disk has opposite effects when the time evolution of the 278 temperatures at the wall region is compared to the location at the mid-axis 279 plane. At the wall location, the phase change process becomes faster for 280 higher values of length and height of the inner rotating disk. The dynamic 281 features of liquid fraction are highly affected with varying length as compared 282 to height (Figure 14). The Fr values become lower with higher length of the 283 rotating disk most of the time. Full phase transition shows non-monotonic 284 behavior with varying geometrical parameters of the inner disk. When the 285 length is increased up to L2d=0.3L, the value of tr is reduced, indicating that 286 the phase change process is relatively fast. Even though the recirculation 287 region becomes larger, due to the increased fluid velocity with rotation, it 288 is lower. However, for disk lengths greater than L2d=0.3L, the value of tr 289 increases, a trend attributed to the movement of the vortex region toward 290 the vertical walls, resulting in an inefficient phase change process. There is 291 a 26% reduction in the value of tr when cases with L2=0.1H and L2=0.3H292 are compared. When the height is increased to H2d=0.1L, the value of 293 tr reduces owing to a reduction of the vortex size in the radial direction. 294

However, further increases of the rotating disk height resulted in longer full phase transition times. The variation in the tr values becomes 21% while the lowest tr value is achieved at H2d=0.1H2.

298 3.2. ANN prediction

Feed-forward ANNs are used for performance predictions and dynamic 299 feature extraction of the PCM-equipped energy system with a rotating disk. 300 As input, five input data is selected. They are Re, Rew, height /length of 301 the rotating disk and time (t in minutes). As network output, the liquid 302 fraction (Lr) is selected. Here, 1877 CFD simulation datasets are generated 303 with 70% used for training and the remainder used for testing and vali-304 dation. As the activation function, a hyperbolic tangent sigmoid function 305 $(f(x) = 1/(1 + e^{(-x)}))$ is used while LM-BP is selected as the learning algo-306 rithm. The number of neurons in the hidden layer is determined according 307 to their performances in the hidden layer; it is chosen to be 15 (Table 2) for 308 training. Figure 16 shows a schematic view of the ANN structure with differ-309 ent layers and network features. ANN performance with 15 neurons and for 310 various data sets (training, validation and testing) is given in Table 3. Lower 311 values of MSE are obtained while R^2 approach 1. Comparison of network 312 performance estimation for various Rew of the inner disk is given in Figure 313 18 (a). Effects of varying the length of the rotating disk on the full phase 314 transition time is given in Figure 18 (b) by using CFD simulation and ANN 315 estimation model. These results show that the ANN model has higher pre-316 diction accuracy when analyzing the effects of using inner rotating cylinder 317 on the phase change dynamics. Variations of the full phase completion time 318 and and time dependent variation of the liquid fraction can be estimated by 319

³²⁰ using the ANN model where time is used as an additional parameter.

321 4. Conclusions

In the present study, the effect of using a rotating inner disk on the performance of a PCM-equipped thermo-fluid system containing a hybrid nanofluid. The following conclusions can be drawn as:

- Introducing rotational surface effects of the inner disk significantly af fected the flow features. There is a 21% reduction in the phase tran sition time at Re=100 when the inner disk is rotating at the highest
 speed as compared to the motionless cases.
- Vortex formation occurs within the system at higher rotational speeds due to the resultant changes in the phase change process. At higher Rew, the phase change process is accelerated with up to a 16 % reduction in the transition time
- Separated flow zones occupying the PCM region are affected by the size of the rotating inner disk.
- The dynamic properties of the liquid fraction are influenced more by varying the length than by varying the height of the inner disk. When configurations at disk sizes of L2=0.1H and L2=0.3H are compared, there is a 26% reduction in the phase transition time but this drops to 21% when varying the height of the rotating disk.
- Feed-forward ANN modeling with 15 neurons in the hidden layer is shown to provide fast and accurate estimation results.

In future work, this study can be extended to include different thermal boundary conditions, different PCMs, geometric modifications in the main reactor and the PCM region, effects of HT fluid inlet temperature variations and using various types of hybrid nanofluids. This should increase the applicability of the present results.

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Property	Value
Melting temperature $(T_m, {}^{\circ}C)$	60
Latent heat of fusion (L, kJ/kg)	213
Density-solid (ρ , kg/m ³)	861
Density-liquid (ρ , kg/m ³)	778
Thermal conductivity-solid (k, $W/m^{o}C$)	0.40
Thermal conductivity-fluid (k, $W/m^{o}C$)	0.15
Specific heat-solid (C_p , J/kg ^o C)	1850
Specific heat-fluid (C_p , J/kg ^o C)	2384

Table 1: Thermo-physical properties of PCM

Number of neurons	MSE -Training	\mathbf{R}^2 -Training
10	6.806×10^{-4}	0.9973
15	1.102×10^{-4}	0.9995
25	2.767×10^{-4}	0.9986

Table 2: Network performance dependence on the neuron in the hidden layer

Data Type	Number of samples	MSE	\mathbb{R}^2
Training	1313	1.102×10^{-4}	0.9995
Validation	282	1.287×10^{-4}	0.9994
Testing	282	1.467×10^{-4}	0.9994

Table 3: ANN results for training, testing and validation using 15 neurons in the hidden layer

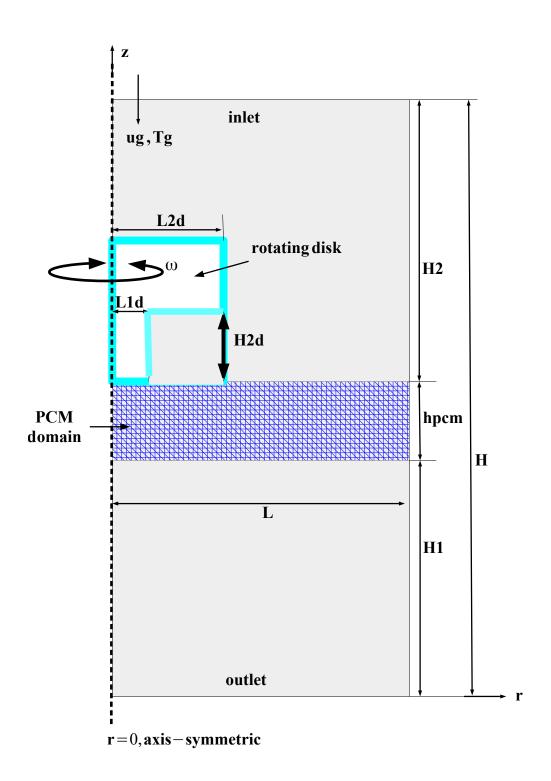


Figure 1: PCM equipped cylinder with a rotating inner disk

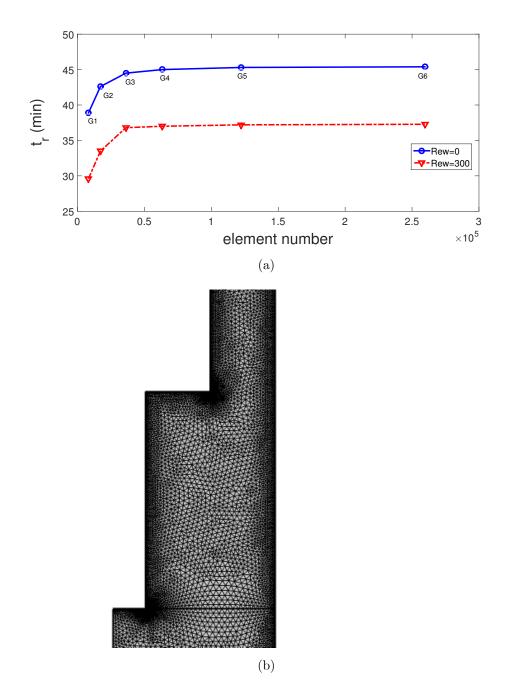


Figure 2: Numerical results for grid independence at two different values of Rew (a) (Re = 200, L1d = 0.1L, L2d = 0.5L, H2d = 0.3H2) and distribution of grid (b)

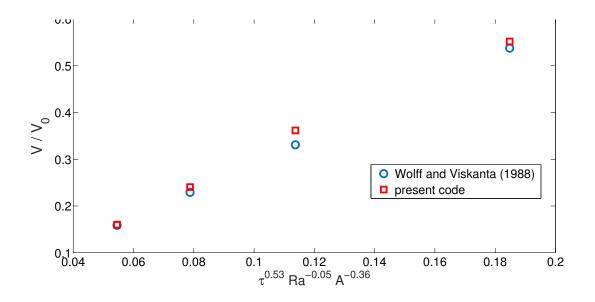


Figure 3: Comparisons of solidified volume fraction for data available in Ref. [41]

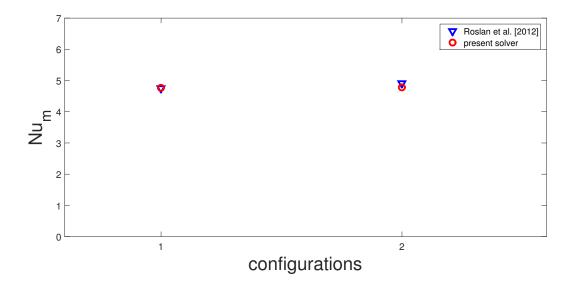


Figure 4: Average Nu comparison in a cavity with an inner rotating cylinder (configuration 1: R=0.1, $\Omega = 500$ and configuration 2: R=0.2, $\Omega = 1000$, Reference values in [42] were used.)

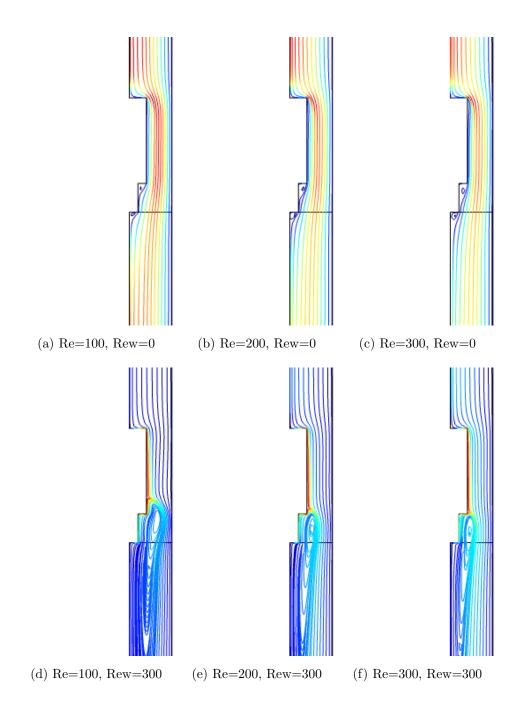


Figure 5: Influence of Re on the flow patterns at two different values of Rew (L1d = 0.1L, L2d = 0.5L, H2d = 0.3H2)

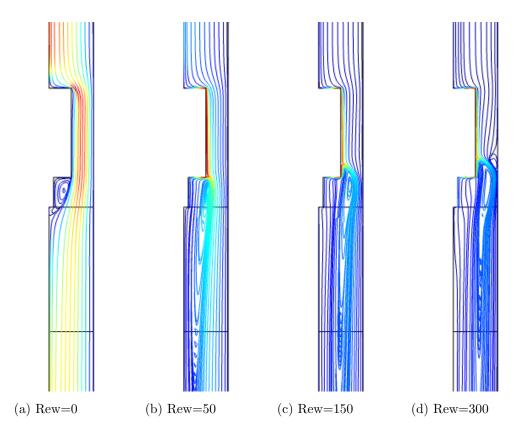


Figure 6: Effects of Rew on the flow patterns (Re = 200, L1d = 0.1L, L2d = 0.5L, H2d = 0.3H2)

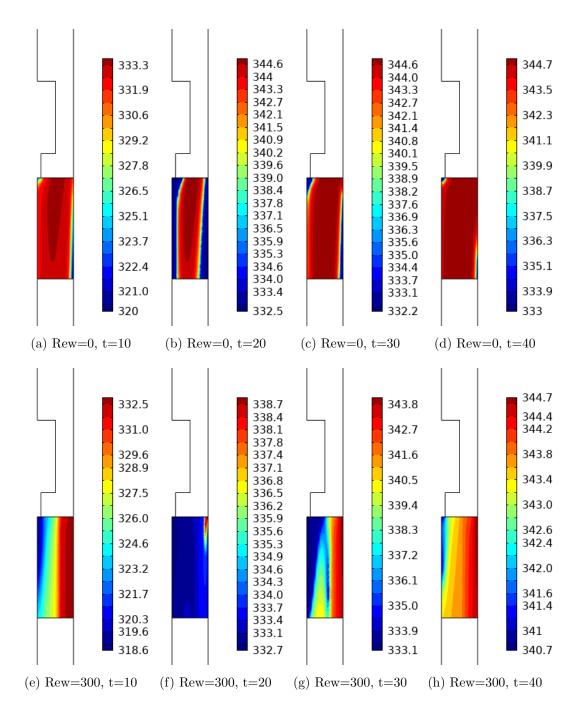


Figure 7: Effects of Rew on Tp at various time instants in minutes (Re = 200, L1d = 0.1L, L2d = 0.5L, H2d = 0.3H2)

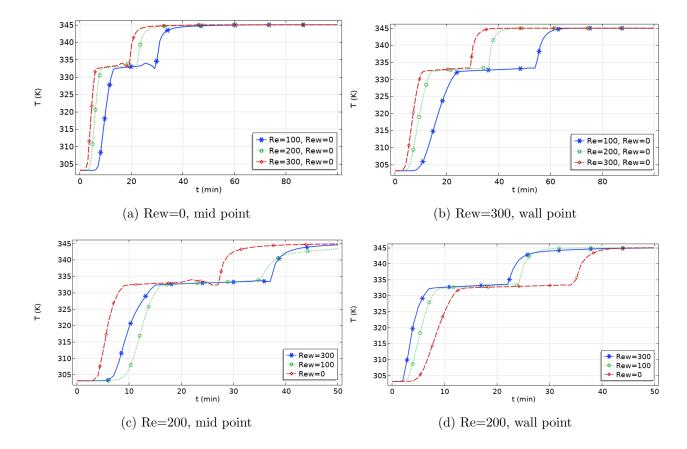
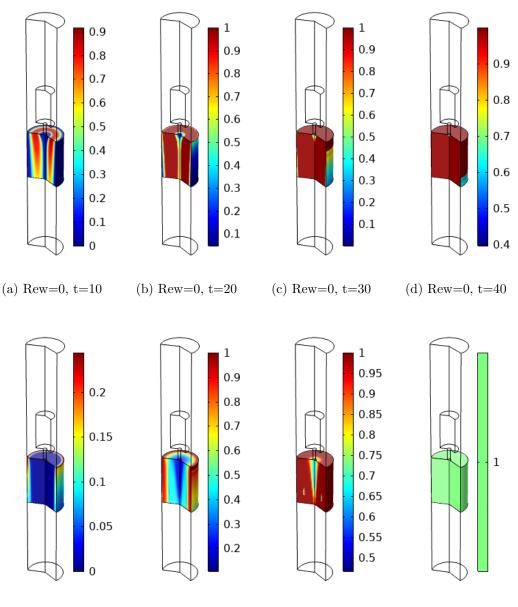


Figure 8: Influence of Re and Rew on the PCM temperature at two different points (Re = 200, L1d = 0.1L, L2d = 0.5L, H2d = 0.3H2)



(e) Rew=300, t=10 (f) Rew=300, t=20 (g) Rew=300, t=30 (h) Rew=300, t=40

Figure 9: Effects of Rew on Lf at various time instants in minutes (Re = 200, L1d = 0.1L, L2d = 0.5L, H2d = 0.3H2)

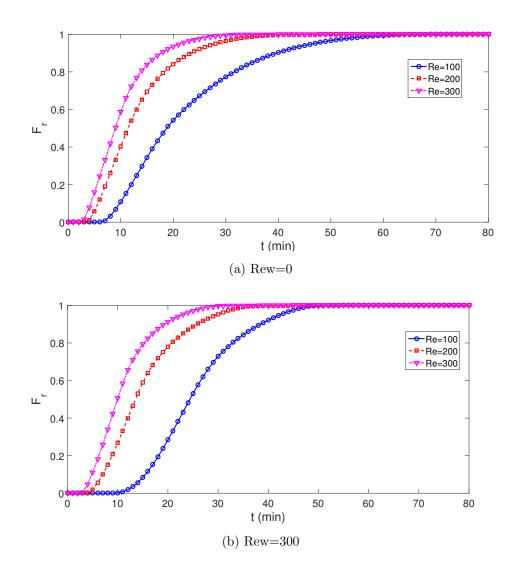


Figure 10: Time evolution of F_r for different Re at two different values of Rew (Re = 200, L1d = 0.1L, L2d = 0.5L, H2d = 0.3H2)

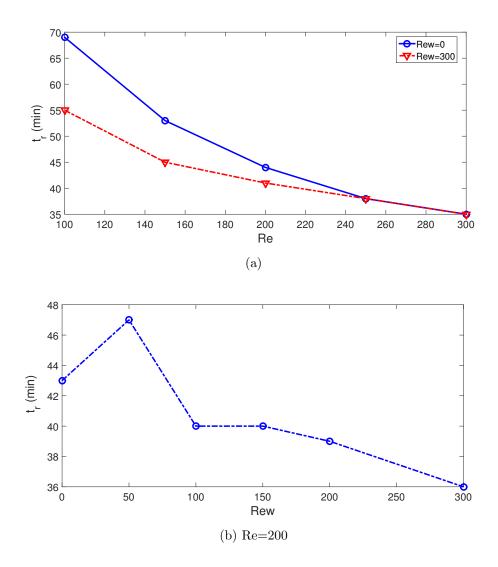


Figure 11: Effects of Re and Rew on the phase transition time (L1d = 0.1L, L2d = 0.5L, H2d = 0.3H2)

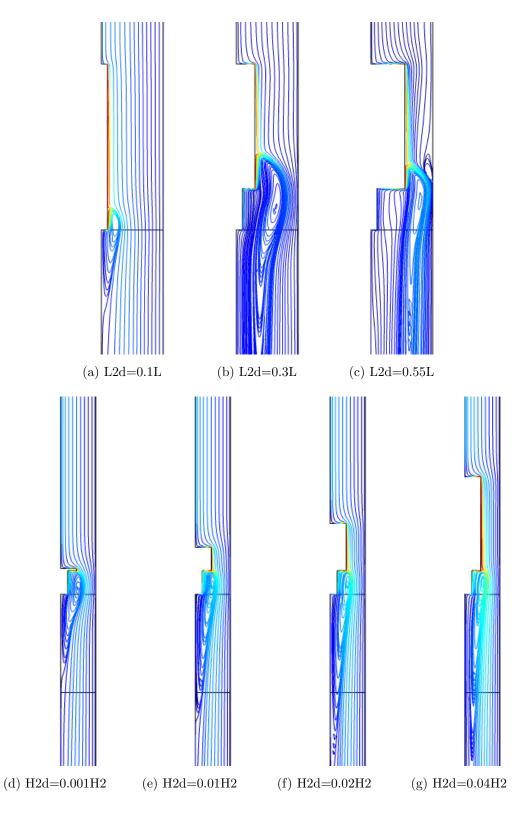


Figure 12: Effects of rotating disk length and height on the flow patterns (Re = 200, Rew=250, L1d=0.1L)

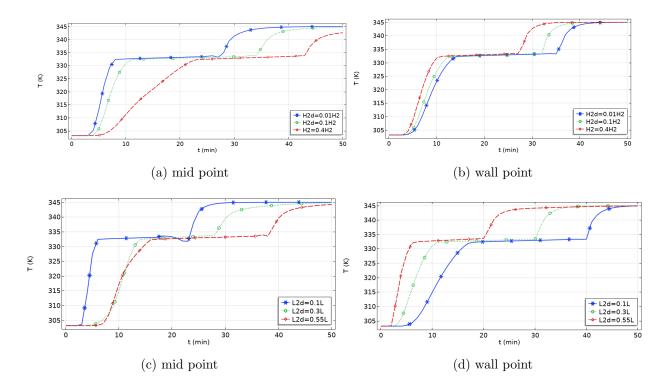


Figure 13: Effects of inner rotating disk height (a-b) and length (c-d) on the PCM temperature variations (Re = 200, Rew = 250, L1d = 0.1L)

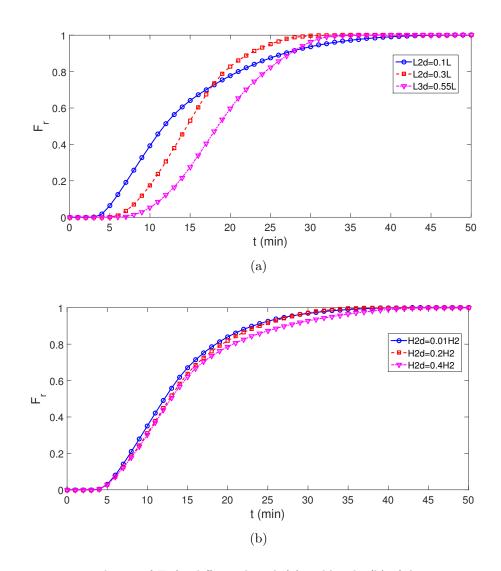


Figure 14: Time evolution of F_r for different length (a) and height (b) of the inner rotating disk (Re = 200, Rew=250, L1d=0.1L)

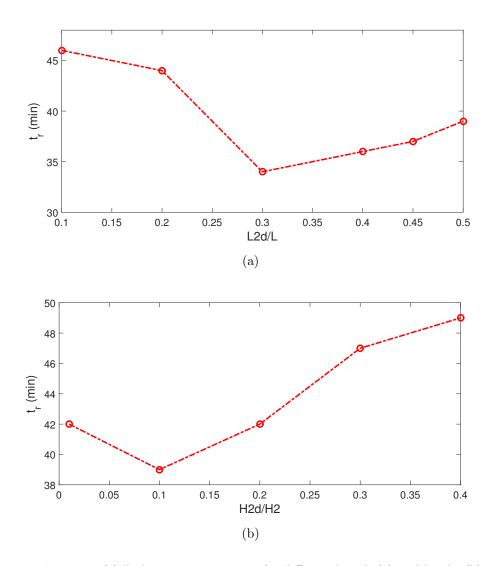


Figure 15: Impacts of full phase transition time for different length (a) and height (b) of the inner rotating disk (Re = 200, Rew = 250, L1d = 0.1L)

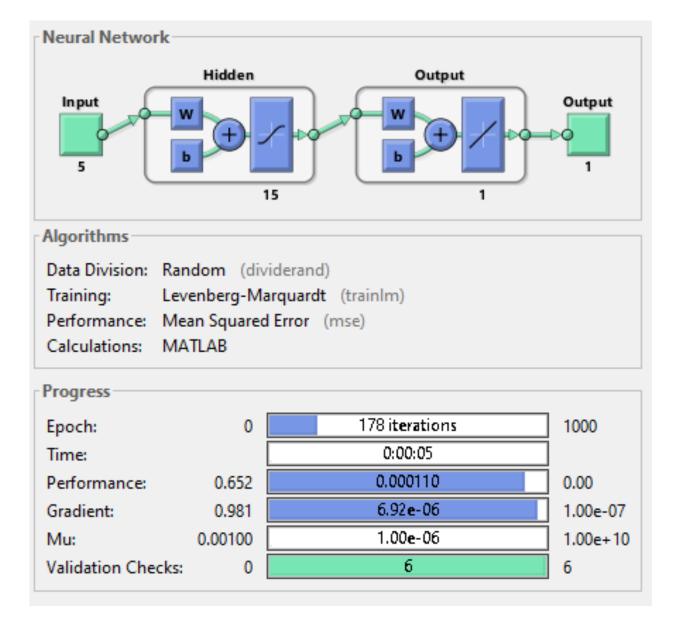


Figure 16: ANN model properties

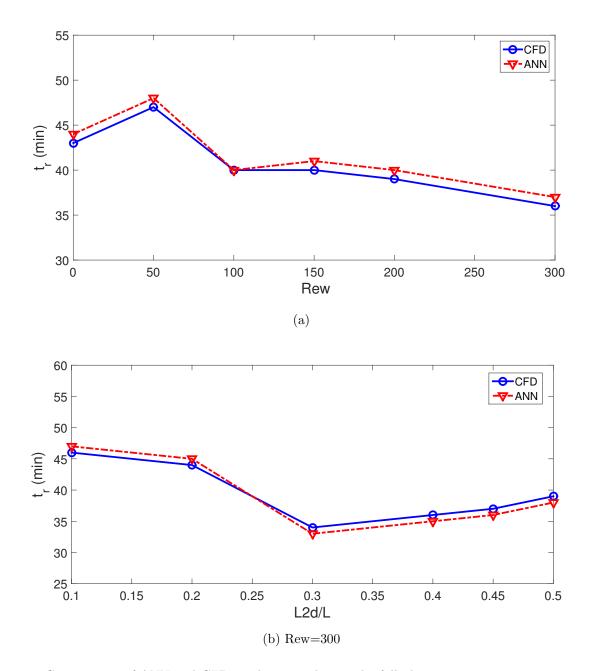


Figure 17: Comparisons of ANN and CFD results in predicting the full phase transition time for varying Rew (b)